# Trajectories of a Quadratic Differential Related to a Particular Algebraic Equation 

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#### Abstract

In this paper, we discuss the existence of a solution interpreted as the Cauchy transform of a signed measure of a particular algebraic quadratic equation of the form $z \mathcal{C}^{2}(z)-P(z) \mathcal{C}(z)+Q(z)=0$ for some polynomials $P(z)$ and $Q(z)$. This issue requires the description of the critical graph of a related quadratic differential in the Riemann sphere $\overline{\mathbb{C}}$. In particular, we discuss the existence of finite critical trajectories of this quadratic differential.


Key words: quantum mechanics, WKB analysis, Cauchy transform, quadratic differentials

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## 1. Introduction

Quadratic differentials have provided an important tool in the asymptotic study of some algebraic equations solutions. In quantum mechanics, trajectories of such quadratic differentials play a crucial role in the WKB analysis.

We consider the algebraic equation

$$
\begin{equation*}
z \mathcal{C}^{2}(z)-z P(z) \mathcal{C}(z)+Q(z)=0, \tag{1.1}
\end{equation*}
$$

where $P(z)$ and $Q(z)$ are two 1-degree real monic polynomials.
Notice that the above technique is the continuity to a series of papers related to the study of complex zeros of hyper-geometric polynomials, for example, $[1,8]$.

In this paper, we discuss the existence of solutions of equation (1.1) as the Cauchy transform of compactly-supported signed measures. In Section 2, we describe the critical graphs of quadratic differentials $-\frac{q(z)}{z} d z^{2}$, where $q$ is a polynomial of degree 3 in the Riemann sphere $\overline{\mathbb{C}}$, precisely, we discuss the existence and the number of its finite critical trajectories. In Section 3, we make the connection between the algebraic equation (1.1) and a particular quadratic differential among those studied in Section 2.

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## 2. A quadratic differential

In the rest of this paper, we denote

$$
\mathbb{C}_{+}=\{z \in \mathbb{C} \mid \Im(z)>0\} ; \quad \mathbb{C}_{-}=\{z \in \mathbb{C} \mid \Im(z)<0\}
$$

Below, we describe the critical graphs of the family of quadratic differentials

$$
\begin{equation*}
\varpi_{a}=-\frac{q(z)}{z} d z^{2}=-\frac{(z-1)(z-a)(z-\bar{a})}{z} d z^{2}, \tag{2.1}
\end{equation*}
$$

where $a \in \mathbb{C}_{+}$, and $q$ is a monic polynomial of degree 3 . We begin our investigation with some immediate observations from the theory of quadratic differentials. For more details, we refer the reader to $[2,6,10]$.

Recall that critical points of a given quadratic differential $-Q(z) d z^{2}$ on the Riemann sphere $\overline{\mathbb{C}}$ are its zeros and poles; the multiplicity of a critical point is its multiplicity in the rational function $Q$ in $\overline{\mathbb{C}}$. Zeros and simple poles are called finite critical points, while poles of order 2 or greater are the infinite ones. All other points of $\overline{\mathbb{C}}$ are called regular points.

Horizontal trajectories (or just trajectories) of the quadratic differential are the zero loci of the equation

$$
-Q(z) d z^{2}>0
$$

or, equivalently,

$$
\begin{equation*}
\Re \int^{z} \sqrt{Q(t)} d t=\text { const. } \tag{2.2}
\end{equation*}
$$

Knowing that if $z(t), t \in \mathbb{R}$, is a horizontal trajectory, we get that the function

$$
t \mapsto \Im \int^{t} \sqrt{Q(z(u))} z^{\prime}(u) d u
$$

is monotone.
The vertical (or orthogonal) trajectories are obtained by replacing $\Im$ by $\Re$ in equation (2.2). The horizontal and vertical trajectories produce two pairwise orthogonal foliations of the Riemann sphere $\overline{\mathbb{C}}$.

A trajectory passing through a critical point is called critical. In particular, if it starts and ends at finite critical points, it is called a finite critical trajectory or a short trajectory. Otherwise, we call it an infinite critical trajectory. The closure of the set of the finite and infinite critical trajectories is called the critical graph. A necessary condition for the existence of a short trajectory connecting two finite critical points is the existence of a Jordan arc $\gamma$ connecting them such that

$$
\begin{equation*}
\Re \int_{\gamma} \sqrt{Q(t)} d t=0 . \tag{2.3}
\end{equation*}
$$

However, this condition is not sufficient in general, see counter-examples in [11].
The local structure of such trajectories is as follows:

- At any regular point, horizontal (respectively, vertical) trajectories look locally as simple analytic arcs passing through this point, and through every regular point, a uniquely determined horizontal (respectively, vertical) trajectory passes; these horizontal and vertical trajectories are orthogonal at this point.
- From each zero of multiplicity $r$, there emanate $r+2$ critical trajectories spacing under equal angle $2 \pi /(r+2)$.
- At a simple pole, there emanates exactly one horizontal trajectory.
- At the pole of order $r>2$, there are $r-2$ asymptotic directions (called critical directions) spacing under equal angle $2 \pi /(r-2)$, and a neighborhood $\mathcal{U}$ such that each trajectory entering $\mathcal{U}$ stays in $\mathcal{U}$ and tends to the pole in one of the critical directions (see Figure 2.1).


Fig. 2.1: Structure of the trajectories near a simple zero (left), a simple pole (center) and a pole of order 6 (right).

A very helpful tool that will be used in our investigation is the so-called Teichmüller Lemma (see [10, Theorem 14.1]).

Definition 2.1. A domain in $\overline{\mathbb{C}}$ bounded only by the segments of horizontal and/or vertical trajectories of $\varpi_{a}$ (and their endpoints) is called a $\varpi_{a}$-polygon.

Lemma 2.2 (Teichmüller). Let $\Omega$ be a $\varpi_{a}$-polygon, and $z_{j}$ be the critical points on the boundary $\partial \Omega$ of $\Omega$, and let $\theta_{j}$ be the corresponding interior angles at vertices $z_{j}$, respectively. Then

$$
\begin{equation*}
\sum\left(1-\frac{\left(n_{j}+2\right) \theta_{j}}{2 \pi}\right)=2+\sum m_{i} \tag{2.4}
\end{equation*}
$$

where $n_{j}$ are the multiplicities of $z_{j}$, and $m_{i}$ the multiplicities of critical points inside $\Omega$.

We have the following immediate observations:

- The finite critical points of $\varpi_{a}$ are simple zeros $1, a, \bar{a}$ and a simple pole at 0 .
- With the parametrization $u=1 / z$, we get

$$
\varpi_{a}(u)=\left(-\frac{1}{u^{6}}+\mathcal{O}\left(\frac{1}{u^{5}}\right)\right) d u^{2}, \quad u \rightarrow 0 .
$$

Thus, infinity is an infinite critical point of $\varpi_{a}$, , namely a pole of order 6 .

- $\quad$ Since the quadratic differential $\varpi_{a}$ has two poles, Jenkins Three-pole Theorem (see [10, Theorem 15.2]) asserts that the situation of the so-called recurrent trajectory (whose closure might be dense in some domain in $\mathbb{C}$ ) cannot happen.
- $\quad$ Since $\infty$ is the only infinite critical point of $\varpi_{a}$, any critical trajectory which is not finite approaches $\infty$ following one of the 4 directions:

$$
D_{k}=\left\{z \in \mathbb{C} \left\lvert\, \arg (z)=(2 k+1) \frac{\pi}{4}\right.\right\}, \quad k=0,1,2,3
$$

Similarly, for the orthogonal trajectories at $\infty$, though the critical directions are:

$$
D_{k}^{\perp}=\left\{z \in \mathbb{C} \left\lvert\, \arg (z)=\frac{k \pi}{2}\right.\right\}, \quad k=0,1,2,3
$$

Observe that if two trajectories approach $\infty$ in the same direction $D_{k}$, then there exists a neighborhood $\mathcal{V}$ of $\infty$, in which any orthogonal trajectory which traverses $D_{k}$ in $\mathcal{V}$ necessarily traverses both of these two trajectories.

Lemma 2.3. Two critical trajectories of $\varpi_{a}$ emanating from the same zero cannot diverge to $\infty$ with the same critical direction.

Lemma 2.4. For any $a \in \mathbb{C}_{+}$, condition (2.3) is fulfilled for $\varpi_{a}$ between the pairs of finite critical points $(0,1)$ and $(a, \bar{a})$. More generally, let $\alpha, \beta$, and $\gamma$ be three complex numbers $(\gamma \neq 0)$. If the quadratic differential

$$
\varpi_{q}=-\frac{q(z)}{z} d z^{2}=-\frac{z^{3}+\alpha z^{2}+\beta z+\gamma}{z} d z^{2}
$$

has two short trajectories connecting two distinct pairs of finite critical points, then

$$
\Im\left(\alpha^{2}-4 \beta\right)=0
$$

In order to study the critical graph of $\varpi_{a}$, we introduce the set

$$
\begin{equation*}
\Sigma=\left\{z \in \mathbb{C} \left\lvert\, \Re \int_{0}^{z} \sqrt{\frac{(t-1)(t-z)(t-\bar{z})}{t}} d t=0\right.\right\} \tag{2.5}
\end{equation*}
$$

Obviously, here our main focus is only in the vanishing of the real part regardless of the choice of the branch-cut of the square root in the integrand.

The following statements will be proved in Section 4.
Lemma 2.5. The set $\Sigma$ is symmetric with respect to the real axis, and it is formed by the 3 Jordan arcs:

- the segment $[0,1]$,
- two curves $\Sigma^{ \pm}$emerging from $z=1$ and diverging respectively to infinity in $\mathbb{C}_{ \pm}$.


Fig. 2.2: Approximate plot of the curve $\Sigma$.
We give here the behavior of $\Sigma$ at $z=1$ and at $\infty$.
Lemma 2.6. The following results hold:

$$
\lim _{\substack{z \rightarrow \infty \\ z \in \Sigma \cap \mathbb{C}^{+}}} \arg (z)=\frac{\pi}{2}, \quad \lim _{\substack{z \rightarrow 1 \\ z \in \Sigma \cap \mathbb{C}^{+}}} \arg (z)=\frac{\pi}{3}
$$

From Lemma 2.5, $\Sigma$ splits $\mathbb{C}$ into two connected domains (see Figure 2.2) :

- $\quad \Omega_{1}$ limited by $\Sigma^{ \pm}$and containing $z=2$,
- $\quad \Omega_{2}=\mathbb{C} \backslash\left(\Omega_{1} \cup \Sigma^{ \pm} \cup[0,1]\right)$.

Proposition 2.7. For any complex number $a \in \mathbb{C}_{+}$, the quadratic differential $\varpi_{a} h a s:$

- two short trajectories if $a \in \Omega_{i}, i=1,2$ : the segment $[0,1]$ and another one that connects $a$ and $\bar{a}$ in $\Omega_{i}$ (see Figures 2.3,2.4);
- three short trajectories if $a \in \Sigma^{ \pm}$: the segment $[0,1]$ and two others that connect $z=1$ with $a$ and $\bar{a}$ that are symmetric with respect to the real axis (see Figure 2.5).

Remark 2.8. The case $\varpi=-\frac{\left(z-a_{1}\right)\left(z-a_{2}\right)\left(z-a_{2}\right)}{z} d z^{2}$, with the zeros satisfying $0<a_{1}<a_{2}<a_{3}$, is obvious. The segments [ $0, a_{1}$ ] and $\left[a_{2}, a_{3}\right]$ are the only two short trajectories (see Figure 2.6 (left)).


Fig. 2.3: Critical graphs when $a \in \Omega_{1}$, here $a=1.6+2 i$ (left) and $a=1.8+2 i$ (right).



Fig. 2.4: Critical graphs when $a \in \Omega_{2}$, here $a=0.5+2 i$ (left) and $a=2 i$ (right).


Fig. 2.5: Critical graphs when $a \in \Sigma$, here $a=1.55+2 i$ (left) and $a=2+6.3 i$ (right).

## 3. Connection with the algebraic equation

The Cauchy transform $\mathcal{C}_{\nu}$ of a compactly-supported complex Borel measure $\nu$ is defined in $\mathbb{C} \backslash \operatorname{supp}(\nu)$ by

$$
\mathcal{C}_{\nu}(z)=\int_{\mathbb{C}} \frac{d \nu(t)}{z-t} .
$$

It has the asymptotics

$$
\mathcal{C}_{\nu}(z)=\frac{\nu(\mathbb{C})}{z}+\mathcal{O}\left(z^{-2}\right), z \rightarrow \infty
$$



Fig. 2.6: Critical graphs for $\varpi_{q}$ when $q=(z-1)(z-2)(z-3)$ (left) and $q=$ $(z-1)(z-2)^{2}$ (right).
and the inversion formula (which should be understood in the distributional sense) reads as

$$
\begin{equation*}
\nu=\frac{1}{\pi} \frac{\partial \mathcal{C}_{\nu}}{\partial \bar{z}} \tag{3.1}
\end{equation*}
$$

In particular, the normalized root-counting measure $\nu_{n}=\nu\left(P_{n}\right)$ of a given complex polynomial $P_{n}$ of degree $n$ in $\mathbb{C}$ is defined by

$$
\nu_{n}=\frac{1}{n} \sum_{P_{n}(a)=0} \delta_{a} \quad(\text { each zero is counted with its multiplicity }) .
$$

Its Cauchy transform is

$$
\mathcal{C}_{\nu_{n}}(z)=\int_{\mathbb{C}} \frac{d \nu_{n}(t)}{z-t}=\frac{P_{n}^{\prime}(z)}{n P_{n}(z)}
$$

While going back to the algebraic equation (1.1), we are seeking a solution everywhere in $\mathbb{C}$ as a Cauchy transform $\mathcal{C}_{\nu}$ of a compactly-supported signed measure $\nu$. With the choice of the square root of the discriminant

$$
\Delta(z)=z\left(z P^{2}(z)-4 Q(z)\right)
$$

of the quadratic equation (1.1) with condition

$$
\sqrt{\Delta(z)} \sim z^{2}, \quad z \rightarrow \infty
$$

through an easy check, we have

$$
\mathcal{C}(z)=\frac{z P(z)-\sqrt{\Delta(z)}}{2 z}=\frac{1}{z}+\mathcal{O}\left(z^{-2}\right), \quad z \rightarrow \infty
$$

Relying on the so-called Plemelj-Sokhotsky formula, it is well known that the measure $\nu$ lies in finite critical trajectories of the quadratic differential $-\frac{\Delta(z)}{z^{2}} d z^{2}$. The measure $\nu$ is given explicitly by

$$
d \nu(t)=\frac{1}{2 i \pi} \frac{(\sqrt{\Delta(t)})_{+}}{t} d t
$$

For more details, we refer the reader to $[3,4,7,9]$.
In the sequel, for general 1-degree real monic polynomials $P(z)$ and $Q(z)$, we may assume that

$$
\Delta(z)=z(z-1)(z-a)(z-\bar{a}), \quad a \in \mathbb{C}_{+}
$$

The following lemma gives a sufficient condition for a solution of (1.1) to be the Cauchy transform of some compactly-supported measure in $\mathbb{C}$.

Lemma 3.1 ([5, Chap. II, Theorem 1.2]). Suppose $f \in L_{\text {loc }}^{1}(\mathbb{C})$ and that $f(z) \rightarrow 0$ as $z \rightarrow \infty$, and let $\mu$ be a compactly-supported measure in $\mathbb{C}$ such that

$$
\mu=\frac{1}{\pi} \frac{\partial f}{\partial \bar{z}}
$$

in the sense of distributions. Then $f(z)=\mathcal{C}_{\mu}(z)$ almost everywhere in $\mathbb{C}$.

Now we announce a theorem summarising the main finding of this paper.
Theorem 3.2. For general 1-degree real monic polynomials $P(z)$ and $Q(z)$, algebraic equation (1.1) has always a solution interpreted as a Cauchy transform of a signed measure $\nu$, supported on the short trajectories $[0,1]$ and $\gamma_{a}$ of the quadratic differential $\varpi_{a}$, and is given explicitly by

$$
\begin{equation*}
d \nu(t)=\frac{1}{2 i \pi} \frac{(\sqrt{\Delta(t)})_{+}}{t} d t \tag{3.2}
\end{equation*}
$$

The measure $\nu$ is of density 1 if and only if

$$
\Re a+(\Im a)^{2}+\frac{15}{4}=0
$$

If equality holds, then $\nu$ is negative on $\gamma_{a}$.

## 4. Proofs

Proof of Lemma 2.3. Suppose that $\gamma_{1}$ and $\gamma_{2}$ are two critical trajectories emanating from the zero $z_{j} \in\{a, 1\}$ and diverging to $\infty$ with the same critical direction $D_{k}$. Consider the $\varpi_{a}$-polygon with edges $\gamma_{1}$ and $\gamma_{2}$, and vertices $z_{j}$ and $\infty$. With the notations of Lemma 2.2, we have

$$
\beta_{j}=\left\{\begin{array}{ll}
0 & \text { if } \theta_{j}=2 \pi / 3, \\
-1 & \text { if } \theta_{j}=4 \pi / 3,
\end{array} \quad \beta_{\infty}=1, \quad \sum m_{i} \geq-1\right.
$$

which violates (2.4).
Proof of Lemma 2.4. Since $\frac{q(t)}{t}$ is a real rational function, then

$$
\begin{equation*}
\sqrt{\frac{q(t)}{t}}=\sqrt{\frac{q(\bar{t})}{\bar{t}}}, \quad t \neq 0 \tag{4.1}
\end{equation*}
$$

So, after changing the variable $u=\bar{t}$ in the second integral, we get

$$
\begin{aligned}
\Re\left(\int_{\bar{z}}^{z} \sqrt{\frac{q(t)}{t}} d t\right) & =\Re\left(\int_{1}^{z} \sqrt{\frac{q(t)}{t}} d t-\int_{1}^{\bar{z}} \sqrt{\frac{q(t)}{t}} d t\right) \\
& =\Re\left(\int_{1}^{z} \sqrt{\frac{q(t)}{t}} d t-\int_{1}^{z} \sqrt{\frac{q(t)}{t}} d t\right) \\
& =\Re\left(2 i \Im\left(\int_{1}^{z} \sqrt{\frac{q(t)}{t}} d t\right)\right)=0
\end{aligned}
$$

Let us provide a necessary condition to get two short trajectories joining two different pairs of finite critical points in the general case of the quadratic differential with simple zeros

$$
\varpi_{q}=-\frac{q(z)}{z} d z^{2}=-\frac{z^{3}+\alpha z^{2}+\beta z+\gamma}{z} d z^{2}, \quad \alpha, \beta, \gamma \in \mathbb{C}, \gamma \neq 0
$$

Considering two disjoint oriented Jordan arcs $\gamma_{1}$ and $\gamma_{2}$ connecting two distinct pairs of finite critical points, we define the single-valued function $\sqrt{\frac{q(z)}{z}}$ in $\mathbb{C} \backslash$ $\left(\gamma_{1} \cup \gamma_{2}\right)$ with the asymptotics $\sqrt{\frac{q(z)}{z}} \sim z, z \rightarrow \infty$. For $s \in \gamma_{1} \cup \gamma_{2}$, we denote by $\left(\sqrt{\frac{q(s)}{s}}\right)_{+}$and $\left(\sqrt{\frac{q(s)}{s}}\right)_{-}$the limits from the + and - sides, respectively. (As usual, the + side of an oriented curve lies to the left and the - side lies to the right if one traverses the curve according to its orientation.)

From the Laurent series of $\sqrt{q(z)}$ at $\infty$, we obtain

$$
\sqrt{\frac{q(z)}{z}}=z+\frac{\alpha}{2}-\left(\frac{\alpha^{2}-4 \beta}{8 z}\right)+\mathcal{O}\left(z^{-2}\right)
$$

Therefore, the residue of $\sqrt{\frac{q(z)}{z}}$ at $z=\infty$ is given by

$$
\operatorname{res}_{z=\infty}\left(\sqrt{\frac{q(z)}{z}}\right)=\frac{1}{8}\left(\alpha^{2}-4 \beta\right) .
$$

Let

$$
I=\int_{\gamma_{1}}\left(\sqrt{\frac{q(s)}{s}}\right)_{+} d s+\int_{\gamma_{2}}\left(\sqrt{\frac{q(s)}{s}}\right)_{+} d s
$$

Since

$$
\left(\sqrt{\frac{q(s)}{s}}\right)_{+}=-\left(\sqrt{\frac{q(s)}{s}}\right)_{-}, \quad s \in \gamma_{1} \cup \gamma_{2},
$$

we have

$$
2 I=\int_{\gamma_{1} \cup \gamma_{2}}\left[\left(\sqrt{\frac{q(s)}{s}}\right)_{+}-\left(\sqrt{\frac{q(s)}{s}}\right)_{-}\right] d s=\oint_{\Gamma} \sqrt{\frac{q(z)}{z}} d z
$$

where $\Gamma$ is a closed contour encircling the curves $\gamma_{1}$ and $\gamma_{2}$. After a deformation of the contour, we pick up the residue at $z=\infty$ and get

$$
I=\frac{1}{2} \oint_{\Gamma} \sqrt{\frac{q(z)}{z}} d z= \pm i \pi \mathrm{res}_{t=\infty}\left(\sqrt{\frac{q(z)}{z}}\right)= \pm \frac{\pi i}{8}\left(\alpha^{2}-4 \beta\right) .
$$

A necessary condition is

$$
\Im\left(\alpha^{2}-4 \beta\right)=0,
$$

which is satisfied for $q=(z-1)(z-a)(z-\bar{a})$.
Proof of Lemma 2.5. Obviously, $\Sigma \cap \mathbb{R}=[0,1]$. The observation (4.1) shows that $\Sigma$ is symmetric with respect to the real axis. In order to prove that $\Sigma$ is a curve, we consider the real functions $F$ and $G$ (locally) defined for $(x, y)$ in $\mathbb{C}_{+}$ by the formulas

$$
F(x, y)=\Re\left(\int_{0}^{x} \sqrt{\frac{(u-(x+i y))(u-(x-i y))(u-1)}{u}} d u\right)
$$

$$
\begin{aligned}
& =\Re\left(\int_{0}^{x} \sqrt{\frac{\left((u-x)^{2}+y^{2}\right)(u-1)}{u}} d u\right) \\
G(x, y) & =\Re\left(\int_{x}^{x+i y} \sqrt{\frac{(u-(x+i y))(u-(x-i y))(u-1)}{u}} d u\right) \\
& =-\int_{0}^{1} y^{2} \sqrt{1-t^{2}} \Im \sqrt{1-\frac{1}{x+i t y}} d t .
\end{aligned}
$$

The square roots are chosen with condition $\sqrt{X}>0$ for $X>0$.
Define

$$
\Sigma=\left\{(x, y) \in \mathbb{R}^{2} \mid(F+G)(x, y)=0\right\}
$$

Let us prove that

$$
\Sigma \backslash[0,1] \subset\{z \in \mathbb{C} \mid \Re z>1\}
$$

Indeed, it is straightforward to check that $F(x, y)=0$ if $0 \leq x \leq 1$ and $y>0$, and $F(x, y) \leq 0$ if $x<0$ and $y>0$. On the other hand, taking the argument in $[0,2 \pi[$, for $0<t \leq 1$, we have

$$
\begin{equation*}
0<\arg (x+i t y)<\arg (x-1+i t y)<\pi \tag{4.2}
\end{equation*}
$$

Therefore,

$$
0<\arg \left(1-\frac{1}{x+i t y}\right)<\pi
$$

implying that

$$
\Im \sqrt{1-\frac{1}{x+i t y}}>0
$$

Thus,

$$
G(x, y)<0
$$

Hence,

$$
(F+G)(x, y) \leq 0+G(x, y)<0, \quad x \leq 1, y>0
$$

As a result, $(x, y) \notin \Sigma$.
Then we prove that $\Sigma$ is a curve, subset of

$$
\Pi=\{(x, y) \mid x>1, y>0\}
$$

For $x>1$, we have

$$
\frac{\partial F}{\partial x}(x, y)=\sqrt{\frac{y^{2}(x-1)}{x}}+\int_{1}^{x} \frac{(x-u)(u-1)}{\sqrt{\left((u-x)^{2}+y^{2}\right)(u-1) u}} d t>0
$$

In addition, for $u_{t}=x+i t y, t \in[0,1]$, we have

$$
\frac{\partial G}{\partial x}(x, y)=\frac{\partial}{\partial x}\left[\Re\left(\int_{0}^{1} i y^{2} \sqrt{1-t^{2}} \sqrt{1-\frac{1}{u_{t}}} d t\right)\right]
$$

$$
=-\int_{0}^{1} \frac{y^{2} \sqrt{1-t^{2}}}{2} \Im\left(\frac{1}{u_{t}^{2} \sqrt{1-\frac{1}{u_{t}}}}\right) d t
$$

It suffices to check that

$$
\forall t \in[0,1] \quad \Im\left(\frac{1}{u_{t}^{2} \sqrt{1-\frac{1}{u_{t}}}}\right) \leq 0
$$

which is equivalent to proving that

$$
\forall t \in[0,1] \quad \arg \left(\frac{1}{u_{t}^{2} \sqrt{1-\frac{1}{u_{t}}}}\right) \in[\pi, 2 \pi[
$$

where the argument is taken in $[0,2 \pi[$. From (4.2), for any $t \in[0,1]$, we get

$$
\left.\arg \left(\frac{1}{u_{t}^{2} \sqrt{1-\frac{1}{u_{t}}}}\right)=2 \pi-\left(\frac{3}{2} \arg \left(u_{t}\right)+\frac{1}{2} \arg \left(u_{t}-1\right)\right) \in\right] \pi, 2 \pi[.
$$

We deduce that for any $t \in[0,1]$,

$$
\Im\left(\frac{1}{u_{t}^{2} \sqrt{1-\frac{1}{u_{t}}}}\right) \leq 0
$$

and then

$$
\frac{\partial G}{\partial x}(x, y) \geq 0
$$

We have just shown that

$$
\frac{\partial(F+G)}{\partial x}(x, y) \neq 0, \quad(x, y) \in \Sigma \cap \Pi
$$

Finally, we conclude that the set $\Sigma$ is a curve in $\mathbb{C}$ by applying the Implicit Function Theorem to the function $F+G$.

Proof of Lemma 2.6. Take $z=r e^{i x} \in \Sigma, r>1, x \in\left[0, \frac{\pi}{2}\right]$. After the change of variable $t=s r e^{i x}$, we get

$$
\Re\left(e^{2 i x} \int_{0}^{1} \sqrt{\frac{\left(s-\frac{1}{r} e^{-i x}\right)(s-1)\left(s-e^{-2 i x}\right)}{s}} d s\right)=0
$$

Taking the limits when $r \rightarrow \infty$, we obtain

$$
\begin{equation*}
0=\Re \int_{0}^{1} e^{2 i x} \sqrt{(s-1)\left(s-e^{-2 i x}\right)} \tag{4.3}
\end{equation*}
$$

Trivially, $x \neq 0$. With the change of variable $t=\alpha u+\beta$, where

$$
\beta=\frac{1+e^{-2 i x}}{2}, \quad \alpha=i \frac{1-e^{-2 i x}}{2}
$$

equation (4.3) becomes

$$
\begin{aligned}
0 & =\Re\left(\int_{\cot x}^{i} \sqrt{u^{2}+1} d u\right) \\
& =\Re\left(\int_{\cot x}^{0} \sqrt{u^{2}+1} d u+\int_{0}^{i} \sqrt{u^{2}+1} d u\right)=\Re\left(\int_{\cot x}^{0} \sqrt{u^{2}+1} d u\right)
\end{aligned}
$$

which holds if and only if $x=\frac{\pi}{2}$.
The Laurent series of $\sqrt{\frac{(t-1)(t-z)(t-\bar{z})}{t}}$ as $t \rightarrow 1$ (with the appropriate choice of the branch-cut of the square root) is

$$
\sqrt{\frac{(t-1)(t-z)(t-\bar{z})}{t}}=|z-1| \sqrt{t-1}+o\left((t-1)^{\frac{1}{2}}\right)
$$

We conclude that

$$
0=\lim _{z \rightarrow 1, z \in \Sigma^{+}} \Re \int_{1}^{z} \sqrt{\frac{(t-1)(t-z)(t-\bar{z})}{t}} d t=\frac{2}{3}|z-1| \Re(z-1)^{\frac{3}{2}}
$$

and then

$$
\arg (z-1)^{\frac{3}{2}} \equiv \frac{\pi}{2} \bmod (\pi)
$$

which ends the proof.
Proof of Proposition 2.7. Clearly, the segment $[0,1]$ is always a short trajectory of $\varpi_{a}$. If $a \notin \Sigma$, then, from (2.3), there is no short trajectory connecting $a$ to 0 or 1. By Lemma 2.3, there exist at most two critical trajectories emanating from $a$ and approaching $\infty$ in the upper half-plane $\mathbb{C}^{+}$. Using the symmetry with respect to the real axis, at least one critical trajectory emanating from $a$ meets a critical trajectory emanating from $\bar{a}$ somewhere at $b \in \mathbb{R} \backslash[0,1]$. Since $b$ cannot be a zero of the quadratic differential $\varpi_{a}$, we conclude that these two critical trajectories form a short one.

If $a \in \Sigma$ and there is no short trajectory connecting $a$ to 1 , then there exist two critical trajectories $\gamma_{a}$ and $\gamma_{1}$ emanating respectively from $a$ and 1 and approaching $\infty$ in the same critical direction $D_{k}$. From the behavior of orthogonal trajectories at $\infty$, we can take an orthogonal trajectory $\sigma$ that hits $\gamma_{1}$ and $\gamma_{a}$ respectively in two points $b$ and $c$ (there are infinitely many such orthogonal trajectories $\sigma$ ). We consider a path $\gamma$ connecting 1 and $a$, formed by the part of $\gamma_{1}$ from 1 to $b$, the part of $\sigma$ from $b$ to $c$, and the part of $\gamma_{a}$ from $c$ to $a$. Then, integrating along $\gamma$, we have

$$
\Re \int_{\gamma} \sqrt{\frac{q(t)}{t}} d t=\Re \int_{1}^{b} \sqrt{\frac{q(t)}{t}} d t+\Re \int_{b}^{c} \sqrt{\frac{q(t)}{t}} d t+\Re \int_{c}^{a} \sqrt{\frac{q(t)}{t}} d t
$$

$$
=\Re \int_{b}^{c} \sqrt{\frac{q(t)}{t}} d t \neq 0
$$

which violates the fact that $a \in \Sigma$.
Proof of Theorem 3.2. We consider the case where the discriminant of algebraic equation (1.1) is

$$
\Delta(z)=z(z-1)(z-a)(z-\bar{a})
$$

for some $a \in \mathbb{C}_{+}$. Let us suppose first that $a \notin \Sigma$ and denote by $\gamma_{a}$ the short trajectory joining $\bar{a}$ to $a$. The segment $[0,1]$ and $\gamma_{a}$ are positively oriented respectively from 0 to 1 , and from $\bar{a}$ to $a$. As in the proof of Lemma 2.4, these orientations define the + and - -sides with respect to the curves $[0,1]$ and $\gamma_{a}$. We choose the square root $\sqrt{\Delta(z)}$ in $\mathbb{C} \backslash\left([0,1] \cup \gamma_{a}\right)$ with asymptotics $\sqrt{\Delta(z)} \sim z^{2}$ as $z \rightarrow \infty$.

From the proof of Lemma 2.4, we have

$$
\begin{aligned}
\nu(\mathbb{C}) & =\int_{[0,1] \cup \gamma_{a}} d \nu(t)=\frac{1}{2 i \pi} \int_{[0,1] \cup \gamma_{a}} \frac{(\sqrt{\Delta(t)})_{+}}{t} d t \\
& =\frac{1}{16}\left(\alpha^{2}-4 \beta\right)=-\frac{1}{4} \Re a-\frac{1}{4}(\Im a)^{2}+\frac{1}{16} .
\end{aligned}
$$

We obtain a necessary and sufficient condition on the zero $a=x+i y, x \in \mathbb{R}, y \geq$ 0 , to get $\nu(\mathbb{C})=1$,

$$
\begin{equation*}
\nu(\mathbb{C})=1 \Longleftrightarrow-y^{2}-\frac{15}{4}=x \tag{4.4}
\end{equation*}
$$

Observe that condition (4.4) cannot hold for $a \in \Sigma$. The expression of the measure $\nu$ on $[0,1]$ is

$$
d \nu(t)_{\mid[0,1]}=\frac{1}{2 i \pi} \frac{(\sqrt{\Delta(t)})_{+}}{t} d t=\frac{1}{2 \pi} \frac{\sqrt{t(1-t)(t-a)(t-\bar{a})}}{t} d t
$$

which obviously implies that it is positive in $[0,1]$. In order to prove that $\nu$ is a non-positive measure in $\gamma_{a}$, we consider the function $f(y)$ defined for $y \geq 0$ by

$$
\begin{aligned}
f(y)=\nu([0,1]) & =\frac{1}{2 \pi} \int_{0}^{1} \frac{\sqrt{t(1-t)(t-a)(t-\bar{a})}}{t} d t \\
& =\frac{1}{2 \pi} \int_{0}^{1} \frac{\sqrt{t(1-t)\left(t^{2}+\left(2 y^{2}+\frac{15}{2}\right) t+y^{4}+\frac{17}{2} y^{2}+\frac{225}{16}\right)}}{t} d t
\end{aligned}
$$

An easy study shows that $f(y)$ increases from $f(0)$ to $\lim _{y \rightarrow+\infty} f(y)=+\infty$. By the other hand,

$$
f(0)=\frac{1}{2 \pi} \int_{0}^{1} \frac{\sqrt{t(1-t)\left(t^{2}+\frac{15}{2} t+\frac{225}{16}\right)}}{t} d t=\frac{1}{2 \pi} \int_{0}^{1} \frac{\left(t+\frac{15}{4}\right) \sqrt{t(1-t)}}{t} d t
$$

$$
=\frac{1}{2 \pi}\left(\mathrm{~B}\left(\frac{3}{2}, \frac{3}{2}\right)+\frac{15}{4} \mathrm{~B}\left(\frac{1}{2}, \frac{3}{2}\right)\right)=\frac{8}{\pi} \mathrm{~B}\left(\frac{3}{2}, \frac{3}{2}\right)=1
$$

We conclude for every $a \in \mathbb{C}_{+}$satisfying (4.4) that

$$
\nu([0,1])>1,
$$

and thus the measure $\nu$ cannot be positive on $\gamma_{a}$.

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## References

[1] M.J. Atia, A. Martínez-Finkelshtein, P. Martínez-González, and F. Thabet, Quadratic differentials and asymptotics of Laguerre polynomials with varying complex parameters, J. Math. Anal. Appl. 416 (2014), 52-80.
[2] Y. Baryshnikov, On stokes Sets, in New Developments in singularity theory ("Cambridge, 2000"), NATO Sci. Ser. II Math. Phys. Chem., 21, Kluwer Acad. Publ., Dordrecht, 2001, 65-86.
[3] T. Bergkvist and H. Rullgård, On polynomial eigenfunctions for a class of differential operators, Math. Res. Lett. 9 (2002), 153-171.
[4] R. Bőgvad and B. Shapiro, On motherbody measures and algebraic Cauchy transform, Enseign. Math. 62 (2016), 117-142
[5] J.B. Garnett, Analytic capacity and measure, LNM, 297, Springer-Verlag, 1972.
[6] J.A. Jenkins, Univalent functions and conformal mapping, Ergebnisse der Mathematik und ihrer Grenzgebiete. Neue Folge, 18, Reihe: Moderne Funktionentheorie, Springer-Verlag, Berlin, 1958.
[7] A. Martínez-Finkelshtein and E.A. Rakhmanov, Critical measures, quadratic differentials, and weak limits of zeros of Stieltjes polynomials, Commun. Math. Phys. 302 (2011) 53-111.
[8] A. Martínez-Finkelshtein, P. Martínez-González, and F. Thabet, Trajectories of Quadratic Differentials for Jacobi Polynomials with Complex Parameters, Comput. Methods Funct. Theory 16 (2016), 347-364.
[9] I.E. Pritsker, How to find a measure from its potential, Comput. Methods Funct. Theory 8 (2008), No. 2, 597-614.
[10] K. Strebel, Quadratic Differentials, 5, Ergebnisse der Mathematik und ihrer Grenzgebiete (3), Springer-Verlag, Berlin, 1984.
[11] F. Thabet, On the existence of finite critical trajectories in a family of quadratic differentials, Bull. Aust. Math. Soc. 94 (2016), 80-91.

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## Траєкторії квадратичного диференціала, пов'язаного з деяким алгебраїчним рівнянням <br> Mondher Chouikhi, Faouzi Thabet, Wafaa Karrou, and Mohamed Jalel

## Atia

У цій статті ми обговорюємо існування розв'язку, інтерпретованого як перетворення Коші деякого заряду, алгебраїчного квадратичного рівняння вигляду $z \mathcal{C}^{2}(z)-P(z) \mathcal{C}(z)+Q(z)=0$ для деяких поліномів $P(z)$ та $Q(z)$. Ця проблема потребує опису критичного графу відповідного квадратичного диференціала на сфері Рімана $\mathbb{C}$. Зокрема, ми обговорюємо існування скінченних критичних траєкторій цього квадратичного диференціала.

Ключові слова: квантова механіка, аналіз WKB, перетворення Коші, квадратичні диференціали


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